

Now that we have these energy eigenstates (1) for hydrogen-like systems,

$$|n\ell m\rangle \equiv \psi_{n\ell m}(r, \theta, \phi) = R_{n\ell}(r) Y_{\ell}^m(\theta, \phi)$$

We can look at their time evolution and superposition.

## Time Evolution

B/c  $|n\ell m\rangle$  are energy eigenstates, the time evolution of a given state is simply,

$$|\psi(t)\rangle \equiv \psi_{n\ell m}(r, \theta, \phi, t) = R_{n\ell}(r) Y_{\ell}^m(\theta, \phi) e^{-iE_n t/\hbar}$$

$$\text{where } E_n = -\frac{1}{2n^2} \left( \frac{ze^2}{4\pi\epsilon_0} \right)^2 \frac{m_e}{\hbar^2} \quad n=1, 2, 3, \dots$$

## Superposition States

If  $|\psi(t)\rangle$  is the result of a superposition of energy eigenstates

then we time evolve each state,

(2)

$$|\psi(t)\rangle \equiv \psi(r, \theta, \phi, t)$$

$$= \sum_{n, l, m} C_{n, l, m} R_{n, l}(r) Y_l^m(\theta, \phi) e^{-iE_n t/\hbar}$$

We can determine coeffs by using projection (as usual)

Note  $E_n$  depends on  $n$  only so that we expect degeneracy!

$$C_{n, l, m} = \langle n, l, m | \psi(t=0) \rangle$$

$$= \int_0^\infty r^2 dr \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi R_{n, l}^*(r) Y_l^{m*}(\theta, \phi) \psi(r, \theta, \phi, 0)$$