

The Solution we have sought thus far is, (1)

$$Y(\theta, \phi) = \Theta(\theta) \Phi(\phi)$$

We have found,

$$\Theta(\theta) = (-1)^m \frac{(2l+1)}{2} \frac{(l-m)!}{(l+m)!} P_l^m(\cos\theta), \quad m \geq 0$$

and

$$\Phi(\phi) = \frac{1}{\sqrt{2\pi}} e^{im\phi}$$

The appropriately normalized product gives,

$$Y_l^m(\theta, \phi) = (-1)^{(m+|m|/2)} \sqrt{\frac{(2l+1)(l-|m|)!}{4\pi(l+|m|)!}} P_l^m(\cos\theta) e^{im\phi}$$

where $l=0, 1, 2, 3, \dots$ and $m=-l, -l+1, \dots, l-1, l$

Our sign choice is conventional and gives,

$$Y_l^{-m}(\theta, \phi) = (-1)^m Y_l^{m*}(\theta, \phi)$$

These are the Spherical Harmonics

and they are position representations of angular momentum eigenstates.

$$|lm\rangle \doteq Y_l^m(\theta, \phi)$$

They are tabulated online and in QM books, ⁽²⁾

$$\begin{array}{ccc} l & m & Y_l^m \\ 0 & 0 & Y_0^0 = \frac{1}{\sqrt{4\pi}} \end{array}$$

$$1 \quad 0 \quad Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos\theta$$

$$1 \quad \pm 1 \quad Y_1^{\pm 1} = \mp \sqrt{\frac{3}{8\pi}} \sin\theta e^{\pm i\phi}$$

etc.

Properties of the Spherical Harmonics

① They are orthonormal on the unit sphere

$$\langle l_1, m_1 | l_2, m_2 \rangle = \int_0^{2\pi} \int_0^\pi Y_{l_1}^{m_1*} Y_{l_2}^{m_2} \sin\theta d\theta d\phi = \delta_{l_1, l_2} \delta_{m_1, m_2}$$

let $d\Omega = \sin\theta d\theta d\phi$ (solid angle)



$$\int Y_{l_1}^{m_1*} Y_{l_2}^{m_2} d\Omega = \delta_{l_1, l_2} \delta_{m_1, m_2}$$

② they form a complete basis

given a smooth $\Psi(\theta, \phi)$,

$$\Psi(\theta, \phi) = \sum_{l=0}^{\infty} \sum_{m=-l}^l c_{lm} Y_l^m(\theta, \phi)$$

the Laplace series

Where coeffs c_{lm} are found via
projection,

(3)

$$\langle lm | \Psi \rangle = c_{lm} = \int_0^{2\pi} \int_0^\pi Y_l^{m*}(\theta, \phi) \Psi(\theta, \phi) \sin\theta d\theta d\phi$$

(3) how Y transforms under parity $\vec{r} \rightarrow -\vec{r}$ depends
on l (angular momentum)

$$Y_l^m(\pi - \theta, \phi + \pi) = (-1)^l Y_l^m(\theta, \phi)$$

(4) they are eigenstates of H , L^2 , & L_z

$$H Y_l^m = \frac{\hbar^2}{2I} l(l+1) Y_l^m$$

$$L^2 Y_l^m = l(l+1) \hbar^2 Y_l^m$$

$$L_z Y_l^m = m \hbar Y_l^m$$

(5) they exhibit degeneracy.

eg.

$E = \frac{\hbar^2}{2I} l(l+1)$ is $2l+1$ degenerate
(for each m).

Thus we must sum over all degenerate
states.

$$P_{E_l} = \sum_{m=-l}^l |\langle lm|\Psi\rangle|^2 \quad \text{prob of } E_l$$

(4)

or for L_z all $l \geq m$ are possible

$$P_{L_z = m\hbar} = \sum_{l=m}^{\infty} |\langle lm|\Psi\rangle|^2$$